

# ESTIMATES FOR THE VOLUME OF A LORENTZIAN MANIFOLD

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ABSTRACT. We prove new estimates for the volume of a Lorentzian manifold and show especially that cosmological spacetimes with crushing singularities have finite volume.

## 0. INTRODUCTION

Let  $N$  be a  $(n + 1)$ -dimensional Lorentzian manifold and suppose that  $N$  can be decomposed in the form

$$(0.1) \quad N = N_0 \cup N_- \cup N_+,$$

where  $N_0$  has finite volume and  $N_-$  resp.  $N_+$  represent the critical past resp. future Cauchy developments with not necessarily a priori bounded volume. We assume that  $N_+$  is the future Cauchy development of a Cauchy hypersurface  $M_1$ , and  $N_-$  the past Cauchy development of a hypersurface  $M_2$ , or, more precisely, we assume the existence of a time function  $x^0$ , such that

$$(0.2) \quad \begin{aligned} N_+ &= x^{0^{-1}}([t_1, T_+)), & M_1 &= \{x^0 = t_1\}, \\ N_- &= x^{0^{-1}}((T_-, t_2]), & M_2 &= \{x^0 = t_2\}, \end{aligned}$$

and that the Lorentz metric can be expressed as

$$(0.3) \quad d\bar{s}^2 = e^{2\psi} \{-dx^{0^2} + \sigma_{ij}(x^0, x)dx^i dx^j\},$$

where  $x = (x^i)$  are local coordinates for the space-like hypersurface  $M_1$  if  $N_+$  is considered resp.  $M_2$  in case of  $N_-$ .

The coordinate system  $(x^\alpha)_{0 \leq \alpha \leq n}$  is supposed to be future directed, i.e. the *past* directed unit normal  $(\nu^\alpha)$  of the level sets

$$(0.4) \quad M(t) = \{x^0 = t\}$$

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is of the form

$$(0.5) \quad (\nu^\alpha) = -e^{-\psi}(1, 0, \dots, 0).$$

If we assume the mean curvature of the slices  $M(t)$  with respect to the past directed normal—cf. [5, Section 2] for a more detailed explanation of our conventions—is strictly bounded away from zero, then, the following volume estimates can be proved

**Theorem 0.1.** *Suppose there exists a positive constant  $\epsilon_0$  such that*

$$(0.6) \quad H(t) \geq \epsilon_0 \quad \forall t_1 \leq t < T_+,$$

and

$$(0.7) \quad H(t) \leq -\epsilon_0 \quad \forall T_- < t \leq t_2,$$

then

$$(0.8) \quad |N_+| \leq \frac{1}{\epsilon_0} |M(t_1)|,$$

and

$$(0.9) \quad |N_-| \leq \frac{1}{\epsilon_0} |M(t_2)|.$$

*These estimates also hold locally, i.e. if  $E_i \subset M(t_i)$ ,  $i = 1, 2$ , are measurable subsets and  $E_1^+, E_2^-$  the corresponding future resp. past directed cylinders, then,*

$$(0.10) \quad |E_1^+| \leq \frac{1}{\epsilon_0} |E_1|,$$

and

$$(0.11) \quad |E_2^-| \leq \frac{1}{\epsilon_0} |E_2|.$$

## 1. PROOF OF THEOREM 0.1

In the following we shall only prove the estimate for  $N_+$ , since the other case  $N_-$  can easily be considered as a future development by reversing the time direction.

Let  $x = x(\xi)$  be an embedding of a space-like hypersurface and  $(\nu^\alpha)$  be the past directed normal. Then, we have the Gauß formula

$$(1.1) \quad x_{ij}^\alpha = h_{ij} \nu^\alpha.$$

where  $(h_{ij})$  is the second fundamental form, and the Weingarten equation

$$(1.2) \quad \nu_i^\alpha = h_i^k x_k^\alpha.$$

We emphasize that covariant derivatives, indicated simply by indices, are always *full* tensors.

The slices  $M(t)$  can be viewed as special embeddings of the form

$$(1.3) \quad x(t) = (t, x^i),$$

where  $(x^i)$  are coordinates of the *initial* slice  $M(t_1)$ . Hence, the slices  $M(t)$  can be considered as the solution of the evolution problem

$$(1.4) \quad \dot{x} = -e^\psi \nu, \quad t_1 \leq t < T_+,$$

with initial hypersurface  $M(t_1)$ , in view of (0.5).

From the equation (1.4) we can immediately derive evolution equations for the geometric quantities  $g_{ij}, h_{ij}, \nu$ , and  $H = g^{ij} h_{ij}$  of  $M(t)$ , cf. e.g. [3, Section 4], where the corresponding evolution equations are derived in Riemannian space.

For our purpose, we are only interested in the evolution equation for the metric, and we deduce

$$(1.5) \quad \dot{g}_{ij} = \langle \dot{x}_i, x_j \rangle + \langle x_i, \dot{x}_j \rangle = -2e^\psi h_{ij},$$

in view of the Weingarten equation.

Let  $g = \det(g_{ij})$ , then,

$$(1.6) \quad \dot{g} = g g^{ij} \dot{g}_{ij} = -2e^\psi H g,$$

and thus, the volume of  $M(t)$ ,  $|M(t)|$ , evolves according to

$$(1.7) \quad \frac{d}{dt} |M(t)| = \int_{M(t_1)} \frac{d}{dt} \sqrt{g} = - \int_{M(t)} e^\psi H,$$

where we shall assume without loss of generality that  $|M(t_1)|$  is finite, otherwise, we replace  $M(t_1)$  by an arbitrary measurable subset of  $M(t_1)$  with finite volume.

Now, let  $T \in [t_1, T_+)$  be arbitrary and denote by  $Q(t_1, T)$  the cylinder

$$(1.8) \quad Q(t_1, T) = \{ (x^0, x) : t_1 \leq x^0 \leq T \},$$

then,

$$(1.9) \quad |Q(t_1, T)| = \int_{t_1}^T \int_M e^\psi,$$

where we omit the volume elements, and where,  $M = M(x^0)$ .

By assumption, the mean curvature  $H$  of the slices is bounded from below by  $\epsilon_0$ , and we conclude further, with the help of (1.7),

$$(1.10) \quad \begin{aligned} |Q(t_1, T)| &\leq \frac{1}{\epsilon_0} \int_{t_1}^T \int_M e^\psi H \\ &= \frac{1}{\epsilon_0} \{|M(t_1)| - |M(T)|\} \\ &\leq \frac{1}{\epsilon_0} |M(t_1)|. \end{aligned}$$

Letting  $T$  tend to  $T_+$  gives the estimate for  $|N_+|$ .

To prove the estimate (0.10), we simply replace  $M(t_1)$  by  $E_1$ .

If we relax the conditions (0.6) and (0.7) to include the case  $\epsilon_0 = 0$ , a volume estimate is still possible.

**Theorem 1.1.** *If the assumptions of Theorem 0.1 are valid with  $\epsilon_0 = 0$ , and if in addition the length of any future directed curve starting from  $M(t_1)$  is bounded by a constant  $\gamma_1$  and the length of any past directed curve starting from  $M(t_2)$  is bounded by a constant  $\gamma_2$ , then,*

$$(1.11) \quad |N_+| \leq \gamma_1 |M(t_1)|$$

and

$$(1.12) \quad |N_-| \leq \gamma_2 |M(t_2)|.$$

*Proof.* As before, we only consider the estimate for  $N_+$ .

From (1.6) we infer that the volume element of the slices  $M(t)$  is decreasing in  $t$ , and hence,

$$(1.13) \quad \sqrt{g(t)} \leq \sqrt{g(t_1)} \quad \forall t_1 \leq t.$$

Furthermore, for fixed  $x \in M(t_1)$  and  $t > t_1$

$$(1.14) \quad \int_{t_1}^t e^\psi \leq \gamma_1$$

because the left-hand side is the length of the future directed curve

$$(1.15) \quad \gamma(\tau) = (\tau, x) \quad t_1 \leq \tau \leq t.$$

Let us now look at the cylinder  $Q(t_1, T)$  as in (1.8) and (1.9). We have

$$(1.16) \quad \begin{aligned} |Q(t_1, T)| &= \int_{t_1}^T \int_{M(t_1)} e^\psi \sqrt{g(t, x)} \leq \int_{t_1}^T \int_{M(t_1)} e^\psi \sqrt{g(t_1, x)} \\ &\leq \gamma_1 \int_{M(t_1)} \sqrt{g(t_1, x)} = \gamma_1 |M(t_1)| \end{aligned}$$

by applying Fubini's theorem and the estimates (1.13) and (1.14).  $\square$

## 2. COSMOLOGICAL SPACETIMES

A cosmological spacetime is a globally hyperbolic Lorentzian manifold  $N$  with compact Cauchy hypersurface  $\mathcal{S}_0$ , that satisfies the timelike convergence condition, i.e.

$$(2.1) \quad \bar{R}_{\alpha\beta} \nu^\alpha \nu^\beta \geq 0 \quad \forall \langle \nu, \nu \rangle = -1.$$

If there exist crushing singularities, see [1] or [2] for a definition, then, we proved in [2] that  $N$  can be foliated by spacelike hypersurfaces  $M(\tau)$  of constant mean curvature  $\tau$ ,  $-\infty < \tau < \infty$ ,

$$(2.2) \quad N = \bigcup_{0 \neq \tau \in \mathbb{R}} M(\tau) \cup \mathcal{C}_0,$$

where  $\mathcal{C}_0$  consists either of a single maximal slice or of a whole continuum of maximal slices in which case the metric is stationary in  $\mathcal{C}_0$ . But in any case  $\mathcal{C}_0$  is a compact subset of  $N$ .

In the complement of  $\mathcal{C}_0$  the mean curvature function  $\tau$  is a regular function with non-vanishing gradient that can be used as a new time function, cf. [4] for a simple proof.

Thus, the Lorentz metric can be expressed in Gaussian coordinates  $(x^\alpha)$  with  $x^0 = \tau$  as in (0.3). We choose arbitrary  $\tau_2 < 0 < \tau_1$  and define

$$(2.3) \quad \begin{aligned} N_0 &= \{(\tau, x) : \tau_2 \leq \tau \leq \tau_1\}, \\ N_- &= \{(\tau, x) : -\infty < \tau \leq \tau_2\}, \\ N_+ &= \{(\tau, x) : \tau_1 \leq \tau < \infty\}. \end{aligned}$$

Then,  $N_0$  is compact, and the volumes of  $N_-, N_+$  can be estimated by

$$(2.4) \quad |N_+| \leq \frac{1}{\tau_1} |M(\tau_1)|,$$

and

$$(2.5) \quad |N_-| \leq \frac{1}{|\tau_2|} |M(\tau_2)|.$$

Hence, we have proved

**Theorem 2.1.** *A cosmological spacetime  $N$  with crushing singularities has finite volume.*

**Remark 2.2.** Let  $N$  be a spacetime with compact Cauchy hypersurface and suppose that a subset  $N_- \subset N$  is foliated by constant mean curvature slices  $M(\tau)$  such that

$$(2.6) \quad N_- = \bigcup_{0 < \tau \leq \tau_2} M(\tau)$$

and suppose furthermore, that  $x^0 = \tau$  is a time function—which will be the case if the timelike convergence condition is satisfied—so that the metric can be represented in Gaussian coordinates  $(x^\alpha)$  with  $x^0 = \tau$ .

Consider the cylinder  $Q(\tau, \tau_2) = \{\tau \leq x^0 \leq \tau_2\}$  for some fixed  $\tau$ . Then,

$$(2.7) \quad |Q(\tau, \tau_2)| = \int_\tau^{\tau_2} \int_M e^\psi = \int_\tau^{\tau_2} H^{-1} \int_M H e^\psi,$$

and we obtain in view of (1.7)

$$(2.8) \quad \tau_2^{-1} \{|M(\tau)| - |M(\tau_2)|\} \leq |Q(\tau, \tau_2)|,$$

and conclude further

$$(2.9) \quad \lim_{\tau \rightarrow 0} |M(\tau)| \leq \tau_2 |N_-| + |M(\tau_2)|,$$

i.e.

$$(2.10) \quad \lim_{\tau \rightarrow 0} |M(\tau)| = \infty \implies |N_-| = \infty.$$

## 3. THE RIEMANNIAN CASE

Suppose that  $N$  is a Riemannian manifold that is decomposed as in (0.1) with metric

$$(3.1) \quad d\bar{s}^2 = e^{2\psi} \{ dx^{0^2} + \sigma_{ij}(x^0, x) dx^i dx^j \}.$$

The Gauß formula and the Weingarten equation for a hypersurface now have the form

$$(3.2) \quad x_{ij}^\alpha = -h_{ij} \nu^\alpha,$$

and

$$(3.3) \quad \nu_i^\alpha = h_i^k x_k^\alpha.$$

As default normal vector—if such a choice is possible—we choose the outward normal, which, in case of the coordinate slices  $M(t) = \{x^0 = t\}$  is given by

$$(3.4) \quad (\nu^\alpha) = e^{-\psi} (1, 0, \dots, 0).$$

Thus, the coordinate slices are solutions of the evolution problem

$$(3.5) \quad \dot{x} = e^\psi \nu,$$

and, therefore,

$$(3.6) \quad \dot{g}_{ij} = 2e^\psi h_{ij},$$

i.e. we have the opposite sign compared to the Lorentzian case leading to

$$(3.7) \quad \frac{d}{dt} |M(t)| = \int_M e^\psi H.$$

The arguments in Section 1 now yield

**Theorem 3.1.** (i) *Suppose there exists a positive constant  $\epsilon_0$  such that the mean curvature  $H(t)$  of the slices  $M(t)$  is estimated by*

$$(3.8) \quad H(t) \geq \epsilon_0 \quad \forall t_1 \leq t < T_+,$$

and

$$(3.9) \quad H(t) \leq -\epsilon_0 \quad \forall T_- < t \leq t_2,$$

then

$$(3.10) \quad |N_+| \leq \frac{1}{\epsilon_0} \lim_{t \rightarrow T_+} |M(t)|,$$

and

$$(3.11) \quad |N_-| \leq \frac{1}{\epsilon_0} \lim_{t \rightarrow T_-} |M(t)|.$$

(ii) *On the other hand, if the mean curvature  $H$  is negative in  $N_+$  and positive in  $N_-$ , then, we obtain the same estimates as Theorem 0.1, namely,*

$$(3.12) \quad |N_+| \leq \frac{1}{\epsilon_0} |M(t_1)|,$$

and

$$(3.13) \quad |N_-| \leq \frac{1}{\epsilon_0} |M(t_2)|.$$

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